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ANNI



UNIVERSITÀ
DEGLI STUDI
DI PADOVA

Dipartimento di Scienze Chimiche – Università degli Studi di Padova
prova finale: Laurea Triennale in
Scienza dei Materiali

development and characterization
of
Superconducting Microwave Cavities
for
Dark Matter search

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Supervisor: Dra. Braggio Caterina

28 October 2022 – AA 2021/2022 – Coorte 2019

Contents

- Axion
- Haloscopes and resonant cavities
- QUAX experiment
- Superconductivity and material properties
- Experimental setup
- Experimental evaluation
- Material treatment and deposition

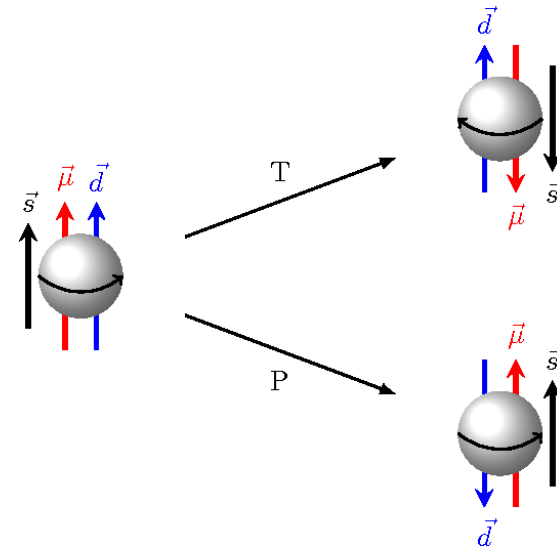
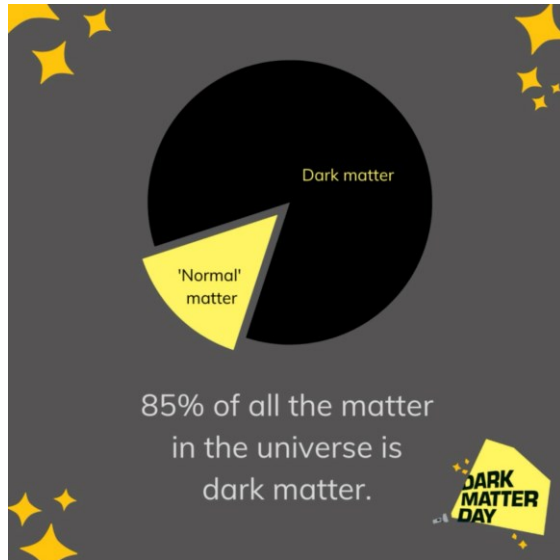
What is an Axion?

non-baryonic Dark Matter candidate

It solves two problems of fundamental physics in one

Weinberg & Wilczek, 1978

STRONG CP PROBLEM



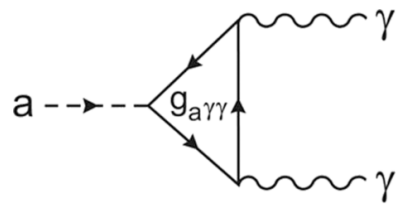
- $m_a \sim 10 - 10^3 \mu eV$
- $\rho_a = 0.45 GeV/cm^3$

Haloscopes

Proposed by P.Sikivie, 1985

Resonant cavity in a strong magnetic field

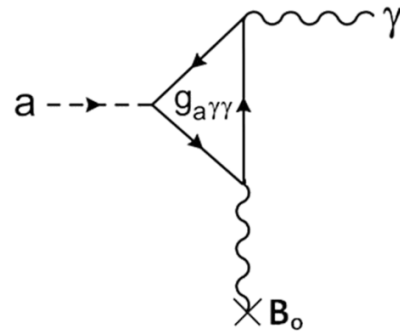
Dilution insert with
microwave cavity
and two magnets



$$\mathcal{L}_{a\gamma\gamma} = g_{a\gamma\gamma} a \vec{E} \cdot \vec{B}$$

$$\tau_{a\gamma\gamma} \approx \left(\frac{10^5 \text{ eV}}{m_a} \right)^5$$

if $\mu\text{eV} \leq m_a \leq \text{meV} \rightarrow \tau_{a\gamma\gamma} \gg \text{age of U}$



$$\mathcal{L}_{a\gamma\gamma} = g_{a\gamma\gamma} a \vec{E} \cdot \vec{B}_0$$

$$\tau_{a\gamma\gamma} \propto B_0^2$$

Inverse Primakoff effect,
conversion of an axion to a photon

Experimental requirements:

- **Multi Tesla field**
SC coils up to 8/10 T

High-Q cavities

Resonant frequency: $\nu_c = \frac{m_a c^2}{h}$

→ $Q \sim 10^6$

Axion signal: $P_a = \left(g_{a\gamma\gamma}^2 \frac{\alpha^2 \hbar^3 c^3 \rho_a}{\pi^2 \Lambda^4} \right) \left(\frac{\beta}{1+\beta} \omega_c \frac{1}{\mu_0} B_0^2 V C_{\text{min}} Q_L \right)$

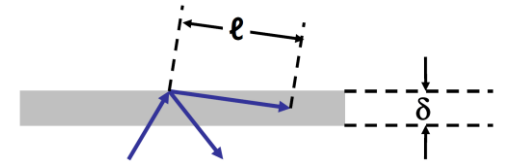
Experimental requirements:

- Multi Tesla field
- Microwave cavity
- $T \sim 10^2 \text{ mK}$
- Ultra high vacuum

Scan rate: $\frac{d\nu}{dt} \propto Q (CV)^2$

$$Q = \frac{\omega_0 U}{P_c} = \frac{G}{R_s}$$





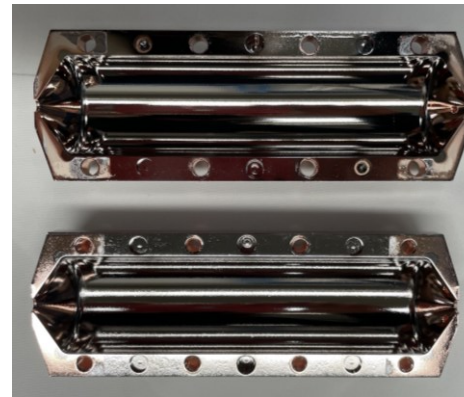
$a\gamma$: Primakoff Haloscope at high frequency ($\sim 6.991 \text{ GHz}$) $\longrightarrow Q_{Cu} \approx 10^5$ (anomalous skin effect)

SC coated Cu cavities

Cu cavity



NbTi coating of about 2-4 μm



Dimensions of the cavity:

- $d_{cil} = 3,2 \text{ cm}$
- $h_{cil} = 12,5 \text{ cm}$
- $h_{con} = 1,1 \text{ cm}$

Superconductivity

Superconducting thermodynamic state

- Zero resistance
- Perfect diamagnetization
- Flux quantization

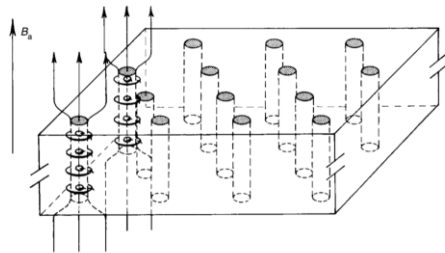
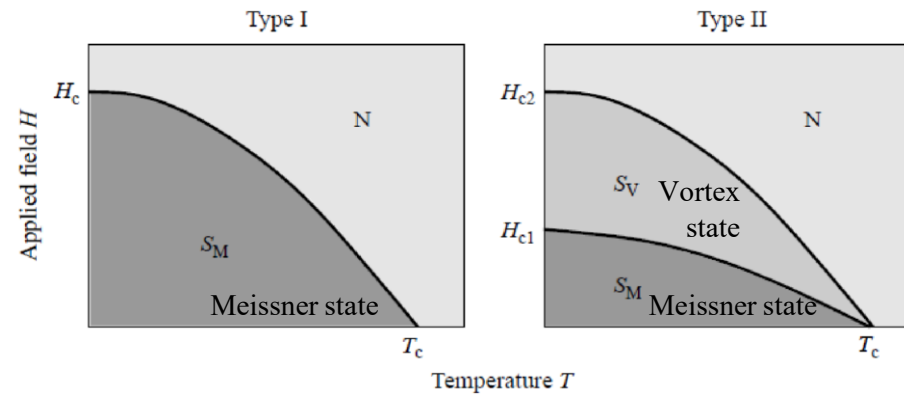
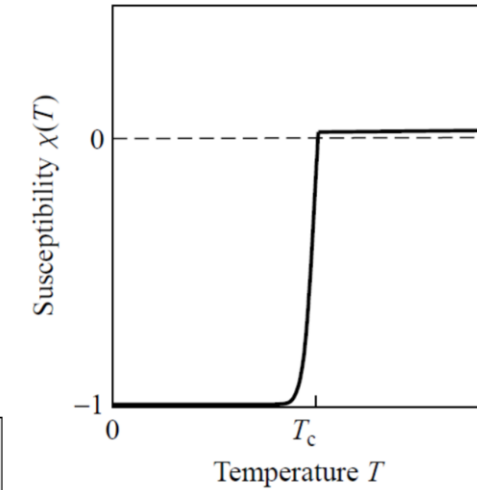
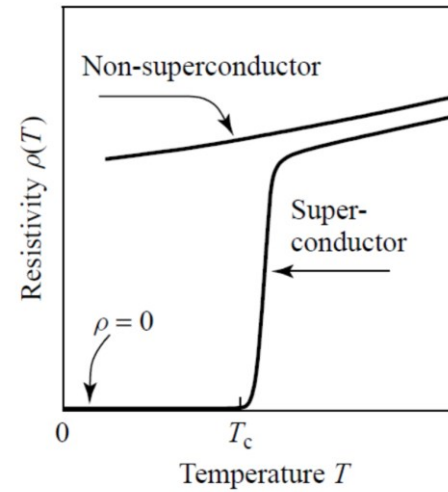


Figure 2.11: Vortices lattices arrangement. B_a is the external applied magnetic field.



NbTi

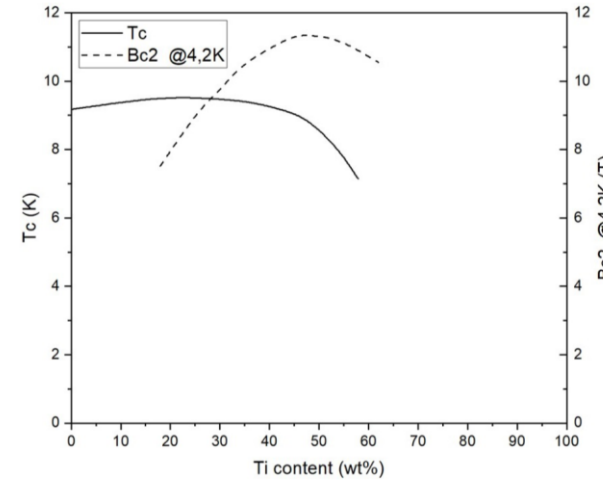
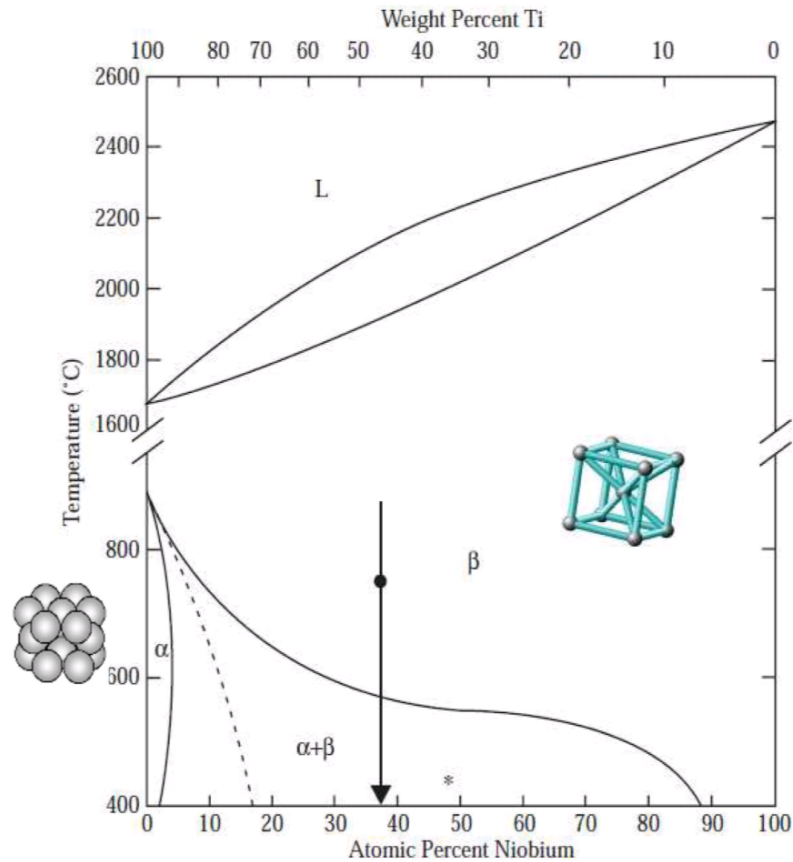


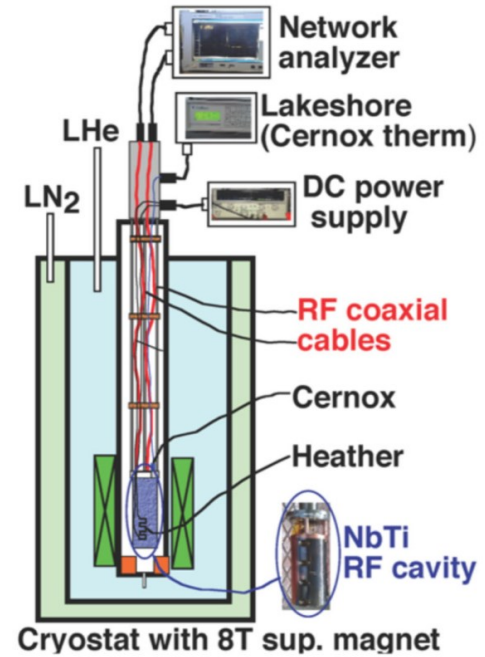
Figure 3.1: Critical temperature and upper critical field of NbTi as a function of the weight percent of Ti. The composition around 50% of Ti offers the highest critical field at the expense of a lower T_c .

$(8 < T_c < 10) K$
 $B_{c2} \leq 12 T$

β : ductile, bcc phase

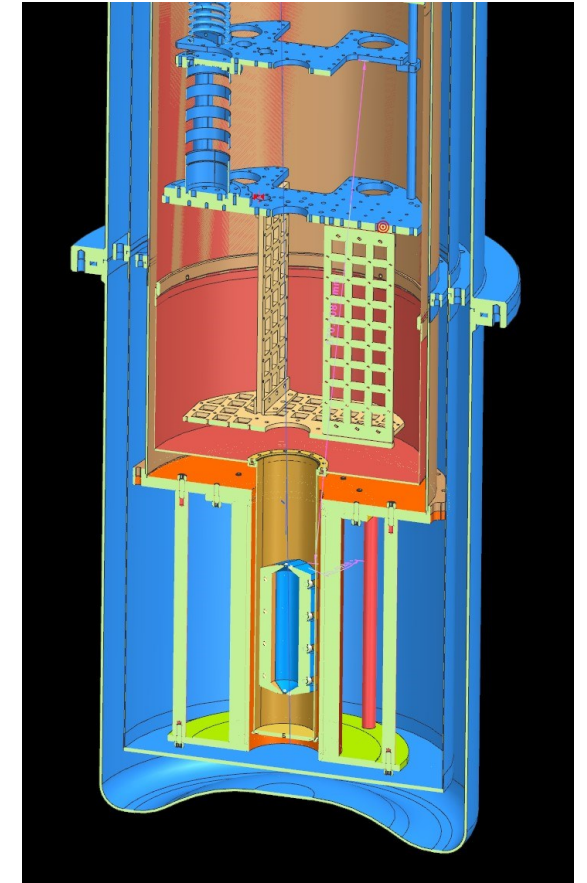
α precipitates: hexagonal non ductile phase

Cavity testing



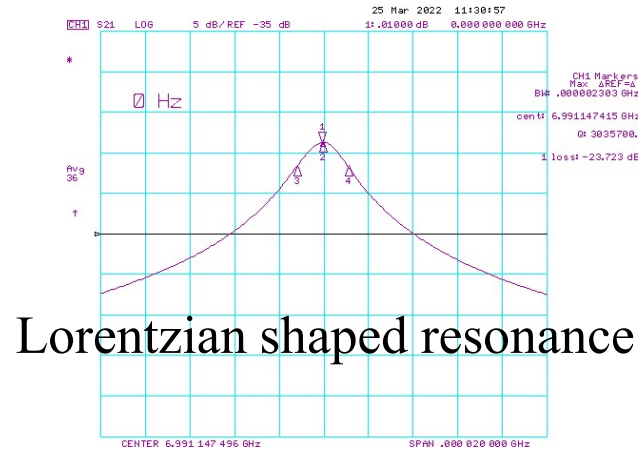
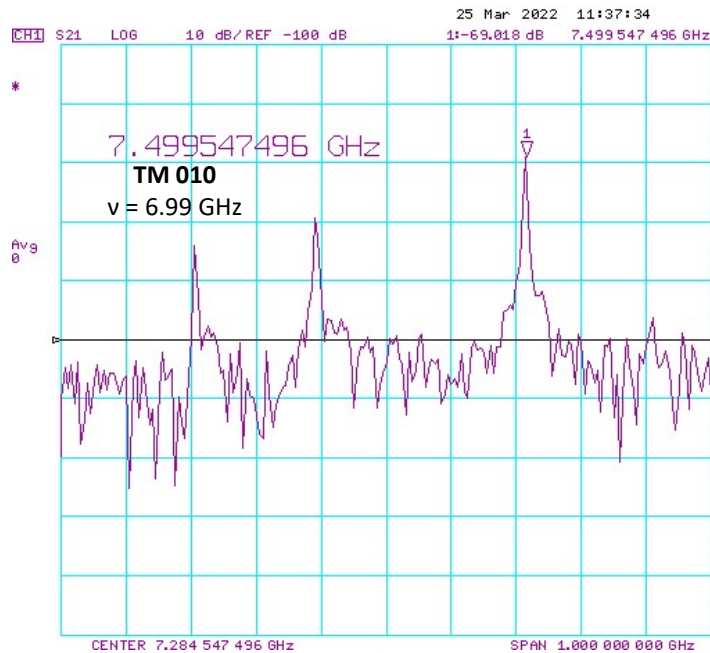
Testing setup:

- VNA
- 2 RF coaxial cables
- He bath
- SC Magnet

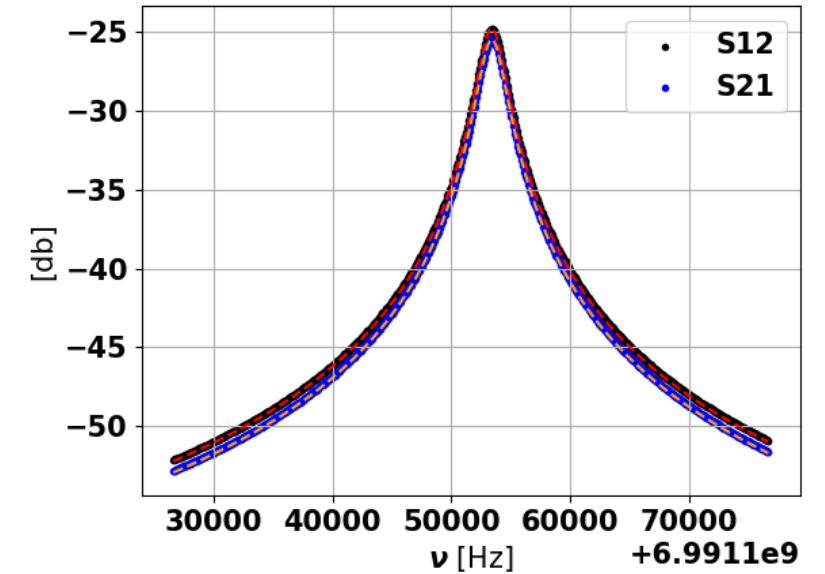


Measurement of Q factor

Transmission channel, S21
Broad spectrum from VNA

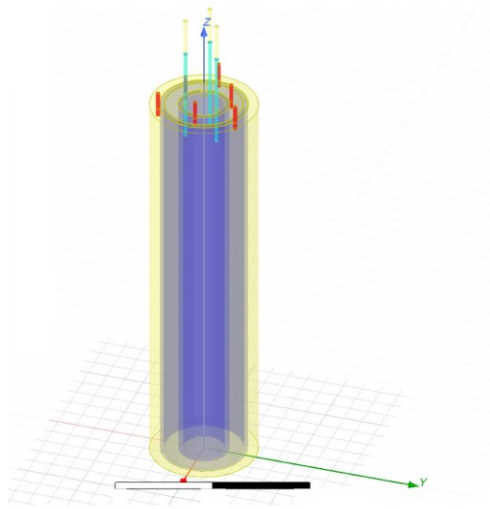


(a) Experimental data and fitting

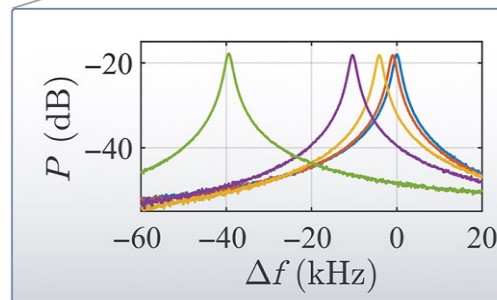
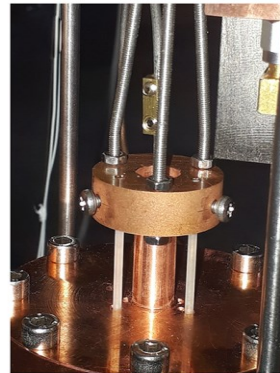
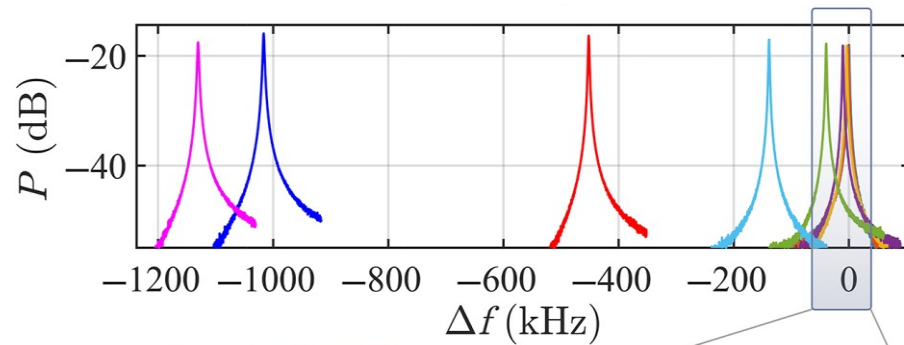


$$Q = \frac{\omega_0 U}{P_c} = \frac{G}{R_s} = \frac{\nu_c}{FWHM}$$

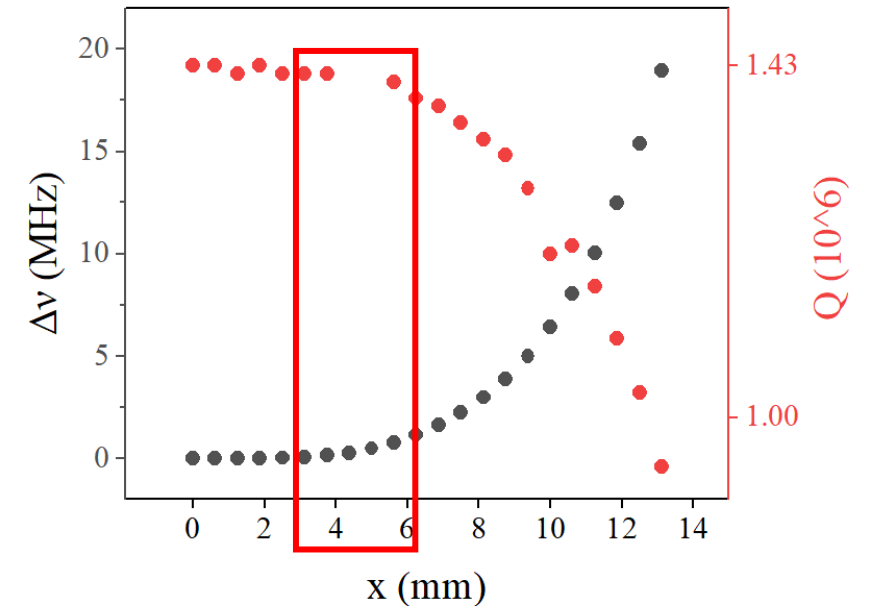
Tuning the cavity frequency



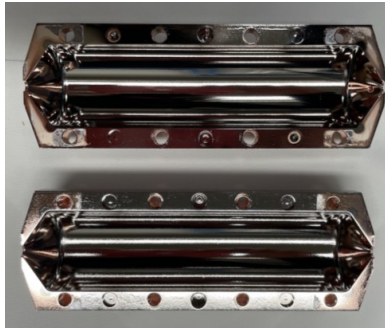
3 adjustable rods of dielectric material
Sapphire (doped $\alpha\text{-Al}_2\text{O}_3$)



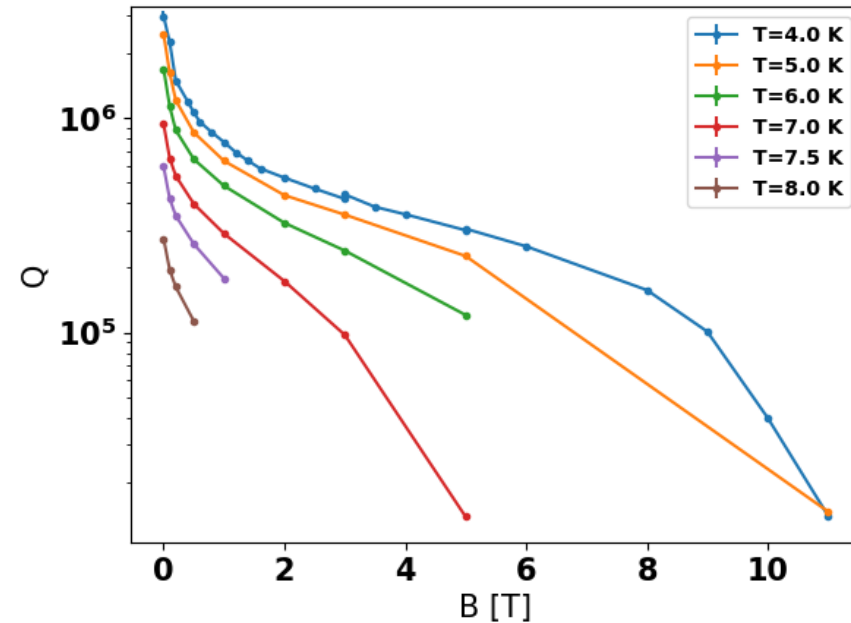
Tuning of ~ 2 MHz starting with the rods 4 mm inside



SC tolerance to intense Magnetic Field



- The higher the B field penetration, the higher the surface resistance, R_s
- Q factor stable up to 3T



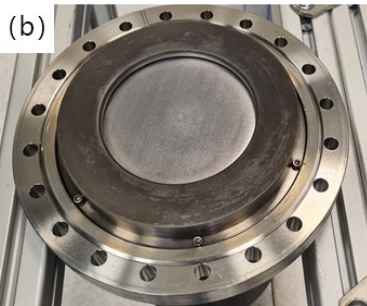
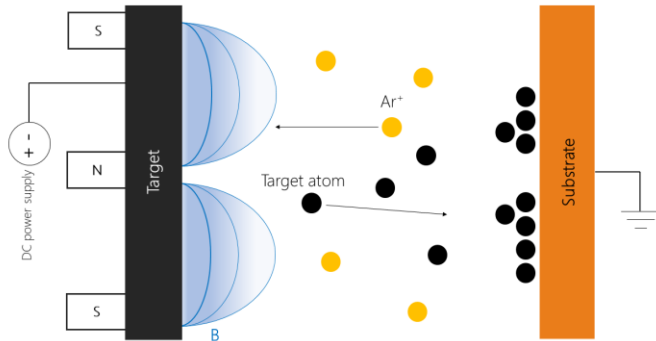
$$Q \approx 5 \times 10^5$$

4.2 K, 3 T, 6.99 GHz

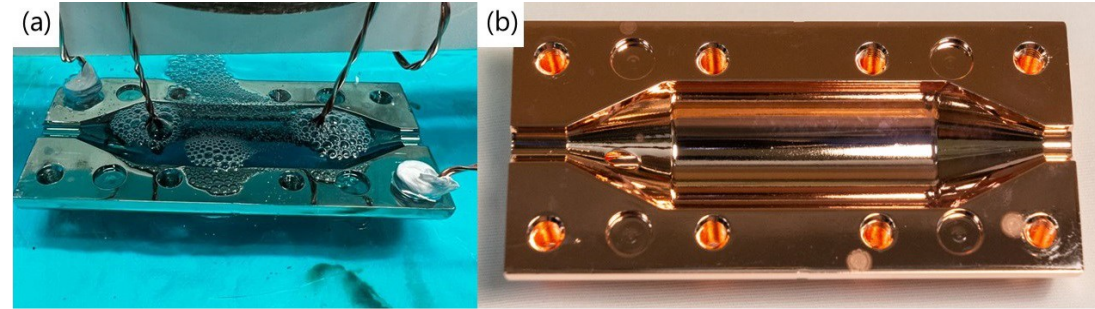
Cavity ready for a test at 3T
(U. Paris-Saclay)

Thin film deposition

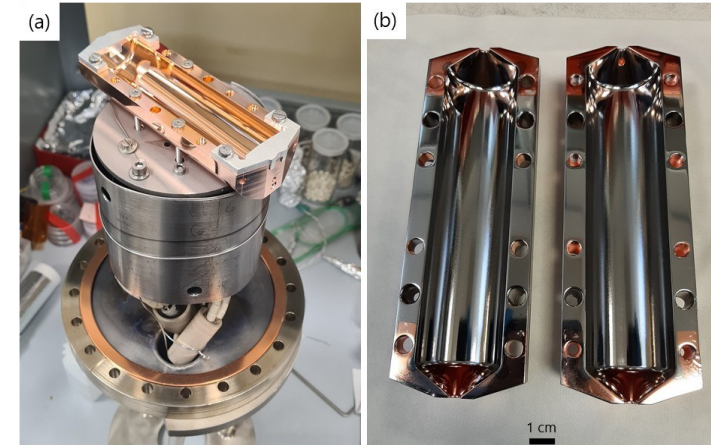
DC Magnetron Sputtering



NbTi target
47 wt%



Non conformal
deposition of
2-4 μm of NbTi on Cu



Surface treatments

- Cleaning and degreasing
- Ultrasonic bath
- EP
- HPR

Table 1: Comparison of the results from this work to several previous studies of superconducting cavities in high magnetic fields for axion research.

Source	Material	f (GHz)	B_a (T)	T (K)	Q_0
This work	Nb ₃ Sn	3.9	6.0	4.2	$(5.3 \pm 0.3) \times 10^5$
[13]	NbTi/Cu	9.08	5	4.2	2.95×10^5
[14]	Nb ₃ Sn	9	8	4.2	6×10^3
[14]	REBCO	9	11.6	4.2	7×10^4
[15]	YBCO	6.93	8.0	4.2	1.5×10^5

Any Questions?

References

- F. Chadha-Day et al., Science Advances, vol. 8, no. 8, Feb. 2022
- Y. K. Semertzidis et al., Science Advances, vol. 8, no. 8, Feb. 2022
- D. Alesini et al., Physical Review D, vol. 99, 101101(R), May 2019
- K. Fossheim et al., «Superconductivity: Physics and Applications», John Wiley & Sons
- R. Kleiner et al., «Superconductivity: Fundamentals and Applications», Wiley-VCH
- A. Salmaso, «Development of NbTi and Nb₃Sn thin films for innovative superconducting axion haloscopes», <https://thesis.unipd.it/handle/20.500.12608/35175>
- H. Padamsee et al., «RF Superconductivity for Accelerators», Cornell University
- C. Mattioli, «Quantum sensing in axion dark matter search», <https://thesis.unipd.it/handle/20.500.12608/34661>
- R. Di Vora et al., Physical Review Applied, vol. 17, no. 5, Feb. 2022
- D. Di Gioacchino et al., IEEE Transactions on Applied Superconductivity, vol. 29, no. 5, pp. 1-5, Aug. 2019
- S. Posen et al., "Measurement of high quality factor superconducting cavities in tesla-scale magnetic fields for dark matter searches." arXiv preprint, arXiv:2201.10733 (2022).

- Part of the images, data and information was gently provided by the QUAX team

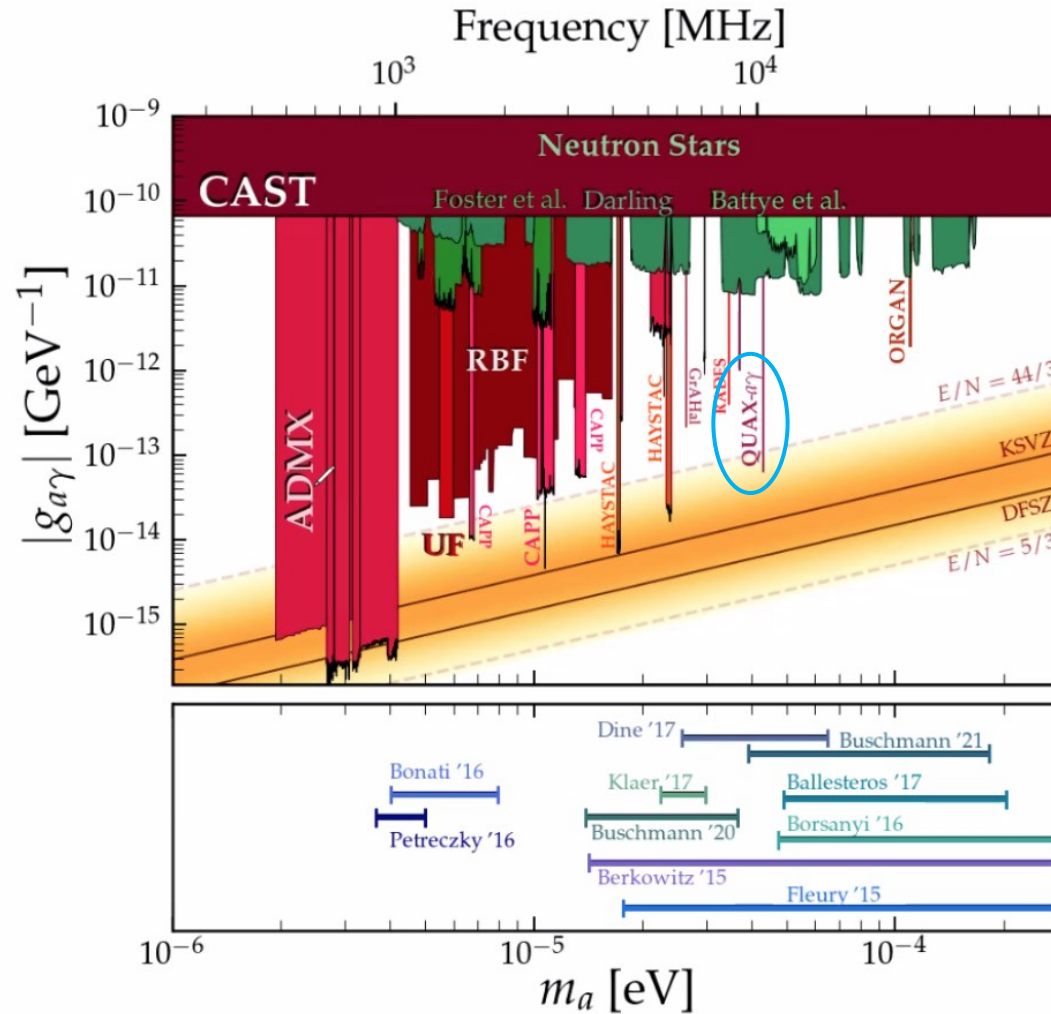
Thank you for your attention!

Exclusion Plot

y-axis: Axion-photon coupling constant, $g_{a\gamma}$
x-axis: Axion mass, m_a

Main theoretical models predict:

$$g_{a\gamma} \propto m_a$$



Axion frequency matched by the resonant cavity

← Target area

Further Cavity characterization

$$Q_L = Q_L(T)$$

$$P_a \propto B^2 Q_0$$

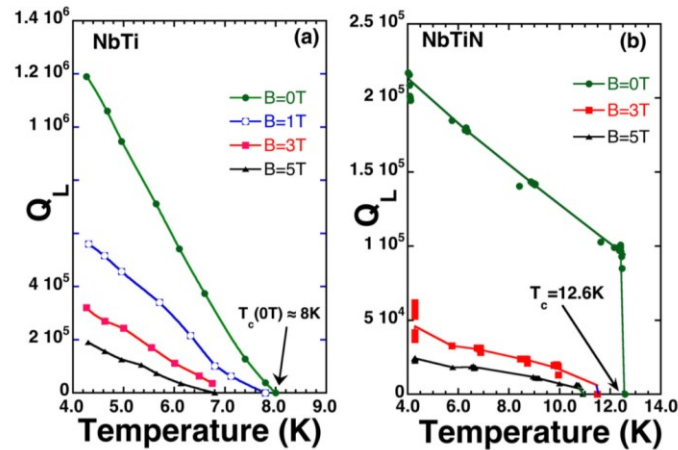
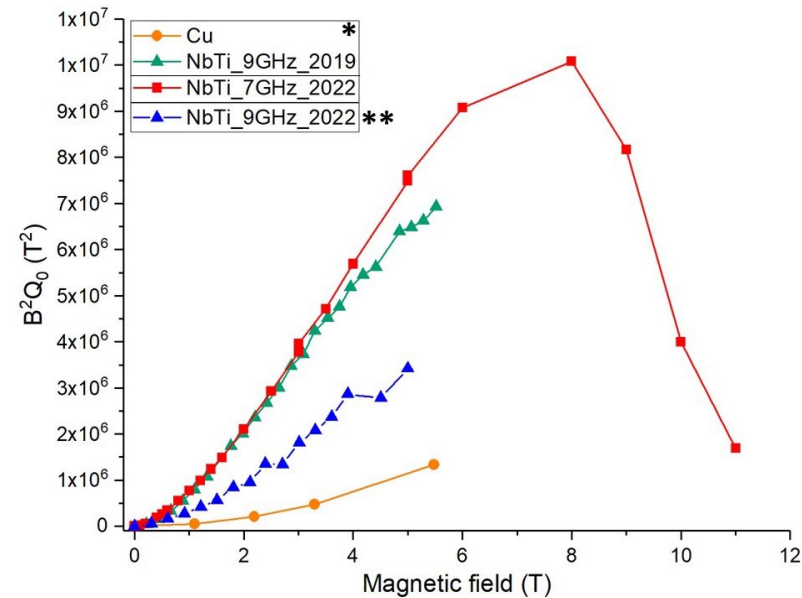


Fig. 2. Q_L versus T at 14.46 GHz (rotated TM_{110}) for (a) NbTi and (b) NbTiN cavities.



Great improvement from copper cavity

$$Q \simeq 3 \times 10^5$$

4.2 K, 4 T, 14.46 GHz

Signal gain with increasing magnetic field

Fluxon motion

IMPEDANCE: $\Delta Z = \Delta R_s + i\Delta X_s$

Gittleman-Rosenblum model advocates the fluxon motion to account for the dissipation

These experimental interpolations allow us to measure with good precision

- London penetration depth, λ
- Depinning frequency, ν_p

in addition to a geometrical parameter for the alignment between fluxons and microwave currents

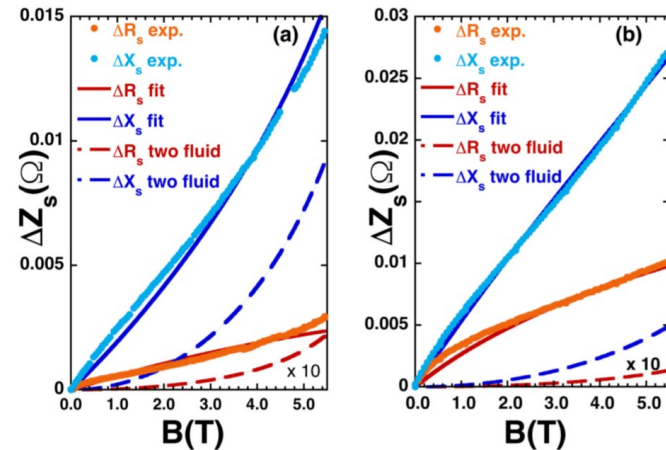


Fig. 5. ΔZ_s vs. B at ~ 4.2 K for the NbTi (a) and NbTiN (b) cavity, mode TM_{110} . Experimental data (full dots), two-fluid only numerical model (dashed lines), two-fluid plus vortex motion numerical model (GR) (continuous line). Parameters values as in the body text.