

Università degli Studi di Padova – Dipartimento di Ingegneria Industriale

Corso di Laurea in Ingegneria Aerospaziale

***Relazione per la prova finale  
«Manifold dynamics in the restricted  
three body problem and applications  
to orbital mechanics»***

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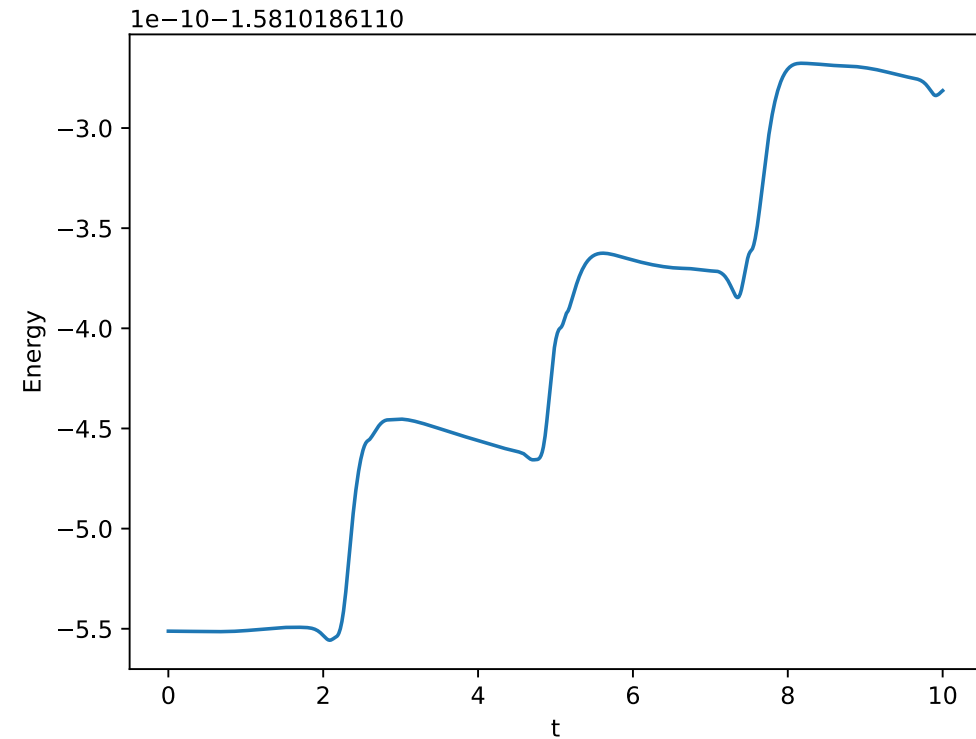
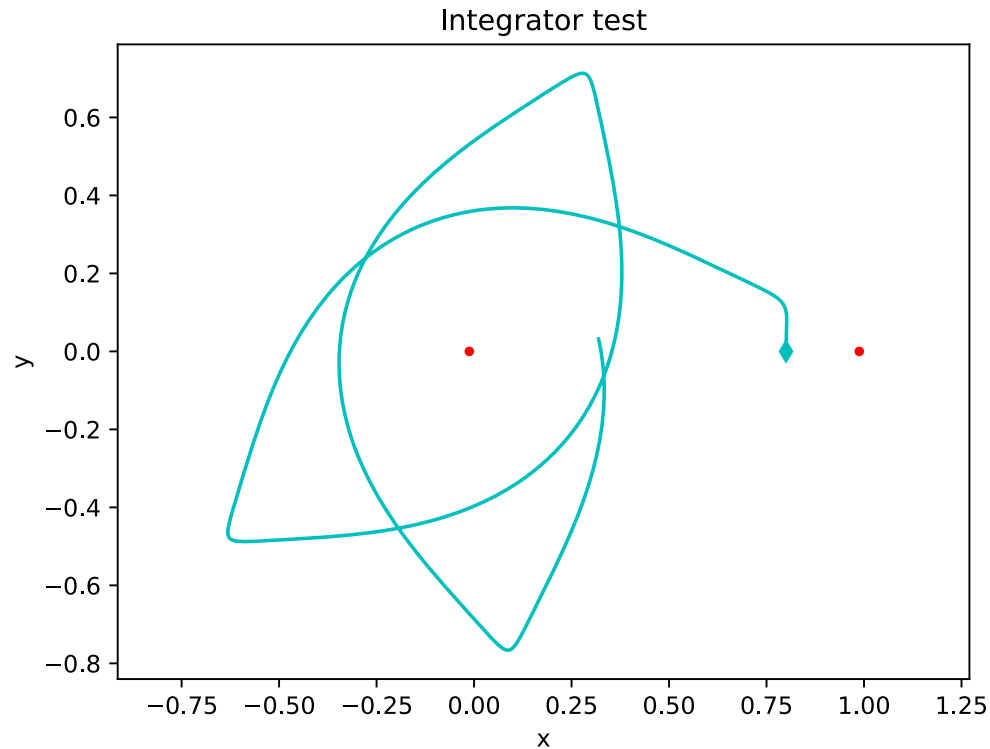
Padova, 16/03/2022

- Simple analysis of the CR3BP
- Collinear Lagrange points
- Linearization of the Hamilton equations
- Lyapunov orbit
- Unstable and stable manifolds of Lyapunov orbit
- Heteroclinic intersection
- Heteroclinic transfer

- The Circular Restricted Three Body Problem, with Hamiltonian

$$H = \frac{1}{2} \left( p_x^2 + p_y^2 + p_z^2 \right) + p_x y - p_y x - \frac{1 - \mu}{\sqrt{(x + \mu)^2 + y^2 + z^2}} - \frac{\mu}{\sqrt{(x + \mu - 1)^2 + y^2 + z^2}}$$

- Check energy conservation to test quality of integration



- Focus on collinear Lagrange points  $L_1, L_2, L_3$

- Equilibrium point in phase space:

$$\underline{c} = (x_c, y_c, p_{x,c}, p_{y,c}) = (x_c, 0, 0, x_c)$$

- Phase state vector:

$$\underline{r} = (x, y, p_x, p_y)$$

- Can be expressed as  $\underline{r} = \underline{c} + \underline{\delta}$

$$x = x_c + \delta x$$

$$y = y_c + \delta y$$

$$p_x = p_{x,c} + \delta p_x$$

$$p_y = p_{y,c} + \delta p_y$$

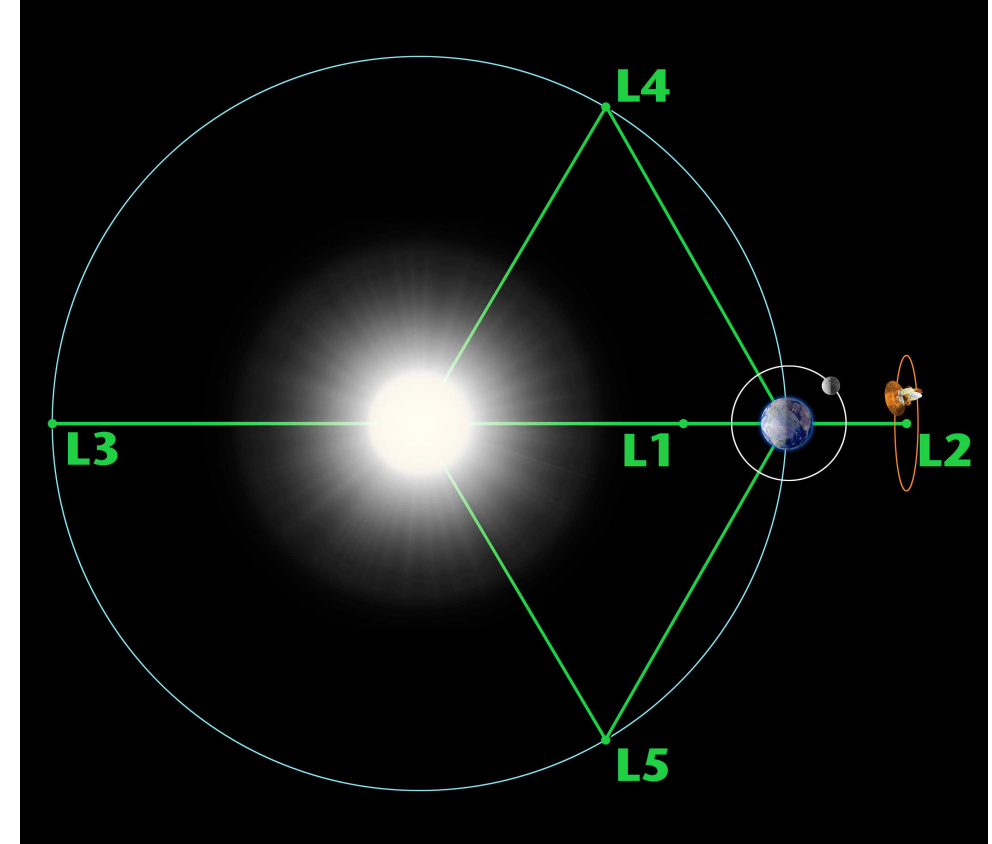
- If  $f(\underline{r})$  are the Hamilton equations, the first order Taylor expansion is:

$$\dot{\underline{r}} = f(\underline{c}) + Df(\underline{c})(\underline{r} - \underline{c}) + \dots$$

$$\dot{\underline{\delta}} = Df(\underline{c})\underline{\delta}$$

- Where  $Df(\underline{r})$  is the Jacobian:

$$\mathbf{J}_{\underline{r}} = \begin{bmatrix} 0 & 1 & 1 & 0 \\ -1 & 0 & 0 & 1 \\ \frac{\partial \dot{p}_x}{\partial x} & \frac{\partial \dot{p}_x}{\partial y} & 0 & 1 \\ \frac{\partial \dot{p}_y}{\partial x} & \frac{\partial \dot{p}_y}{\partial y} & -1 & 0 \end{bmatrix}$$

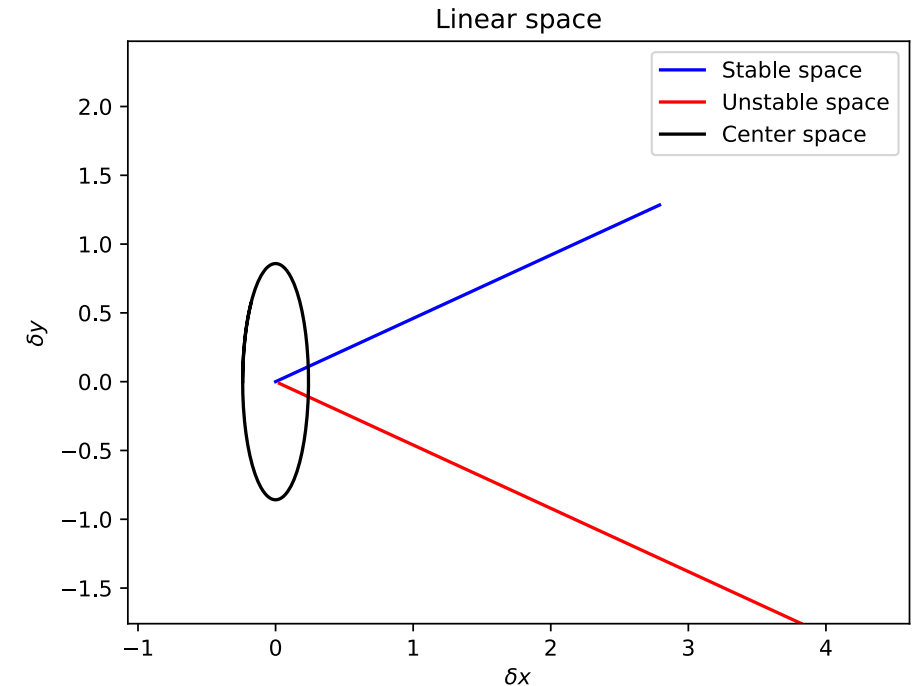


- The eigenvalues of the Jacobian give the linear stability of  $\underline{c}$ :  
 $\lambda_1 = \lambda \quad \lambda_2 = -\lambda \quad \lambda_3 = i\nu \quad \lambda_4 = -i\nu$   
 therefore  $\underline{c}$  is a partially hyperbolic equilibrium point
- The eigenvectors of the Jacobian give the unstable, stable and center spaces, respectively  
 $\mathbf{E}^U = \langle \mathbf{v}_1 \rangle$   
 $\mathbf{E}^S = \langle \mathbf{v}_2 \rangle$   
 $\mathbf{E}^C = \langle \mathbf{v}_3, \mathbf{v}_4 \rangle$
- The local manifold theorem states that, for a system with partially hyperbolic equilibrium, the manifolds exist and are tangent in the equilibrium point to their respective linear spaces
- The stable, unstable and center manifolds are defined as the starting conditions in the phase space  $D$  that:  

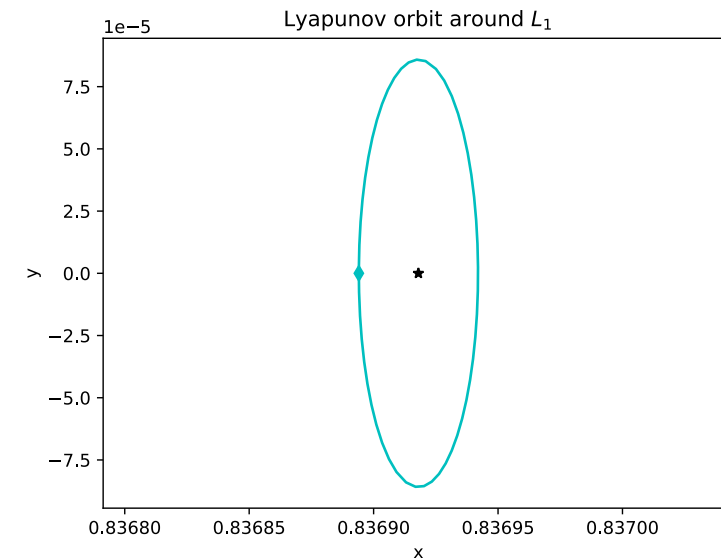
$$W^S = \{ \underline{r} \in D \setminus \underline{c} : \lim_{t \rightarrow +\infty} \phi(t, \underline{r}) = \underline{c} \}$$

$$W^U = \{ \underline{r} \in D \setminus \underline{c} : \lim_{t \rightarrow -\infty} \phi(t, \underline{r}) = \underline{c} \}$$

$$W^C = \{ \underline{r} \in D \setminus \underline{c} : \|\phi(t, \underline{r}) - \underline{c}\| < \epsilon \} \quad \epsilon > 0, t \in \tau$$
- The family of horizontal Lyapunov orbits foliate the center manifold  $W^c$



- Using a starting condition on the center space gives a trajectory that does not make a full orbit, the nonlinear terms dominate the dynamics
- We define the section of dimension 3:  
 $\Sigma = \{ \underline{r} : y = 0 \quad \dot{y} > 0 \}$
- And the section map:  
 $\Sigma \rightarrow \Sigma \quad (x, p_x, p_y) \mapsto (x', p'_x, p'_y) = \psi(x, p_x, p_y)$
- The Lyapunov orbit is a fixed point for the map  $\psi(x, p_x, p_y)$
- The Lyapunov orbit is obtained by finding the root of  
 $F(p_y) = x' - x_0 = \psi_x(x_0, 0, p_y) - x_0$   
using Newton's method



- We can find the energy of the Lyapunov orbit,  $E$
- We define the Poincaré section of dimension 2:  
 $\Pi = \{ \underline{r} : y = 0 \quad \dot{y} > 0, \quad H = E \}$
- And we define the Poincaré map:  
 $\Pi \rightarrow \Pi \quad \underline{s} \mapsto \underline{s}' = \underline{g}(\underline{s})$   
 where  $\underline{s} = (x, p_x)$
- The fixed point  $\underline{p}$  of the Lyapunov orbit is again a fixed point for the Poincaré map
- We can do a first order Taylor expansion around the fixed point:

$$\underline{s}' = \underline{p} + M(\underline{s} - \underline{p}) + \dots$$

$$\underline{\eta}' = M\underline{\eta}$$

where  $\underline{\eta} = \underline{s} - \underline{p}$  and  $M$  is the monodromy matrix:

$$M = \begin{bmatrix} \frac{\partial g_x}{\partial x} & \frac{\partial g_x}{\partial p_x} \\ \frac{\partial g_{p_x}}{\partial x} & \frac{\partial g_{p_x}}{\partial p_x} \end{bmatrix}$$

- The partial derivatives for the monodromy matrix are calculated through finite differences of first order:

$$\frac{\partial g_x}{\partial x} \approx \frac{g_x(x_0 + h, 0) - g_x(x_0 - h, 0)}{2h}$$

$$\frac{\partial g_x}{\partial p_x} \approx \frac{g_x(x_0, h) - g_x(x_0, -h)}{2h}$$

$$\frac{\partial g_{p_x}}{\partial x} \approx \frac{g_{p_x}(x_0 + h, 0) - g_{p_x}(x_0 - h, 0)}{2h}$$

$$\frac{\partial g_{p_x}}{\partial p_x} \approx \frac{g_{p_x}(x_0, h) - g_{p_x}(x_0, -h)}{2h}$$

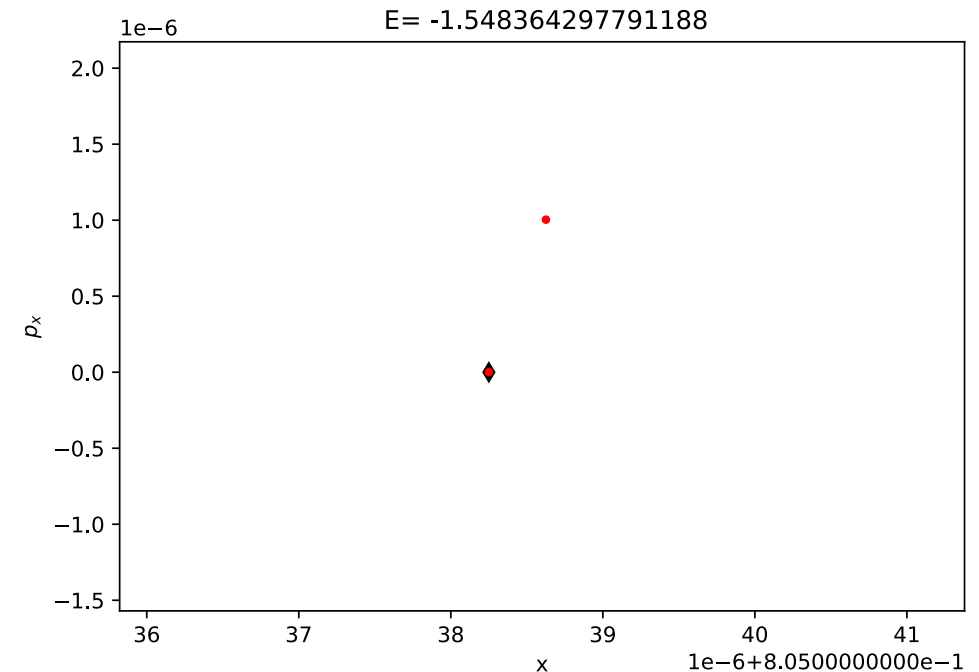
- Check that the monodromy matrix is symplectic

$$\det(M) = 1$$

- The eigenvalues of the monodromy matrix give the linear stability of the fixed point

$$|\rho_1| > 1 \quad |\rho_2| < 1$$

- The respective eigenvectors generate the unstable and stable spaces of the linearized map





- **Definition:** the unstable and stable manifolds of  $\underline{p}$  are the sets:

$$W_{\underline{p}}^U = \{x, p_x \in \Sigma \setminus \underline{p} : \lim_{k \rightarrow +\infty} g^k = \underline{p}\}$$

$$W_{\underline{p}}^S = \{x, p_x \in \Sigma \setminus \underline{p} : \lim_{k \rightarrow -\infty} g^k = \underline{p}\}$$

- **Proposition:** the sets  $W_{\underline{p}}^U$   $W_{\underline{p}}^S$  are invariant under the action of the map  $\underline{g}(\underline{s})$

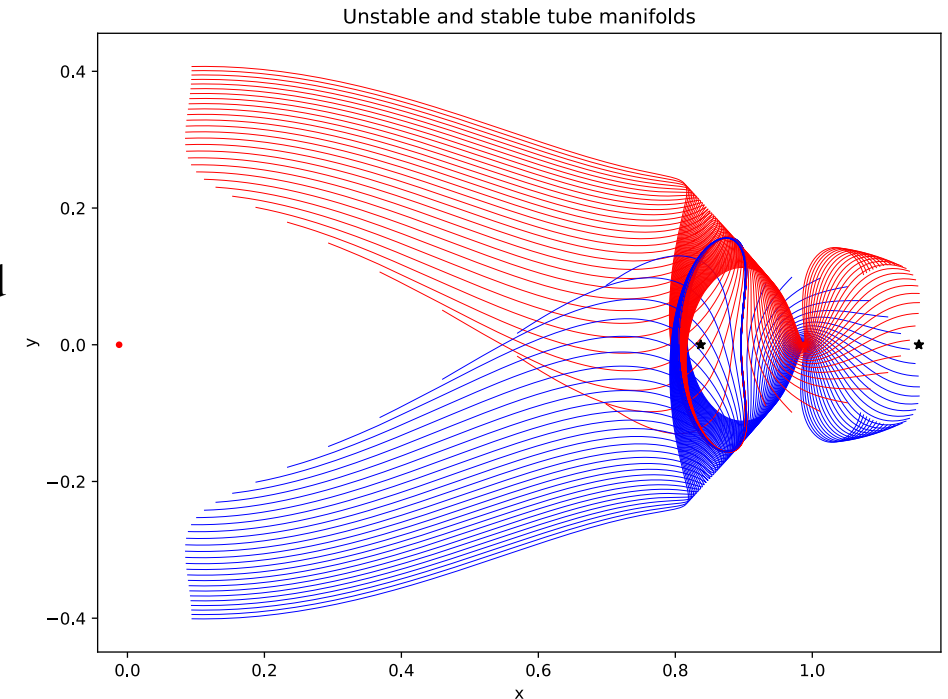
- **Theorem** (Local manifold theorem): The manifolds  $W_{\underline{p}}^U$   $W_{\underline{p}}^S$  exist, and are tangent to the respective invariant manifolds of the fixed point of the linearized map  $\underline{\eta}' = M\underline{\eta}$

- To visualise the stable tube manifold:

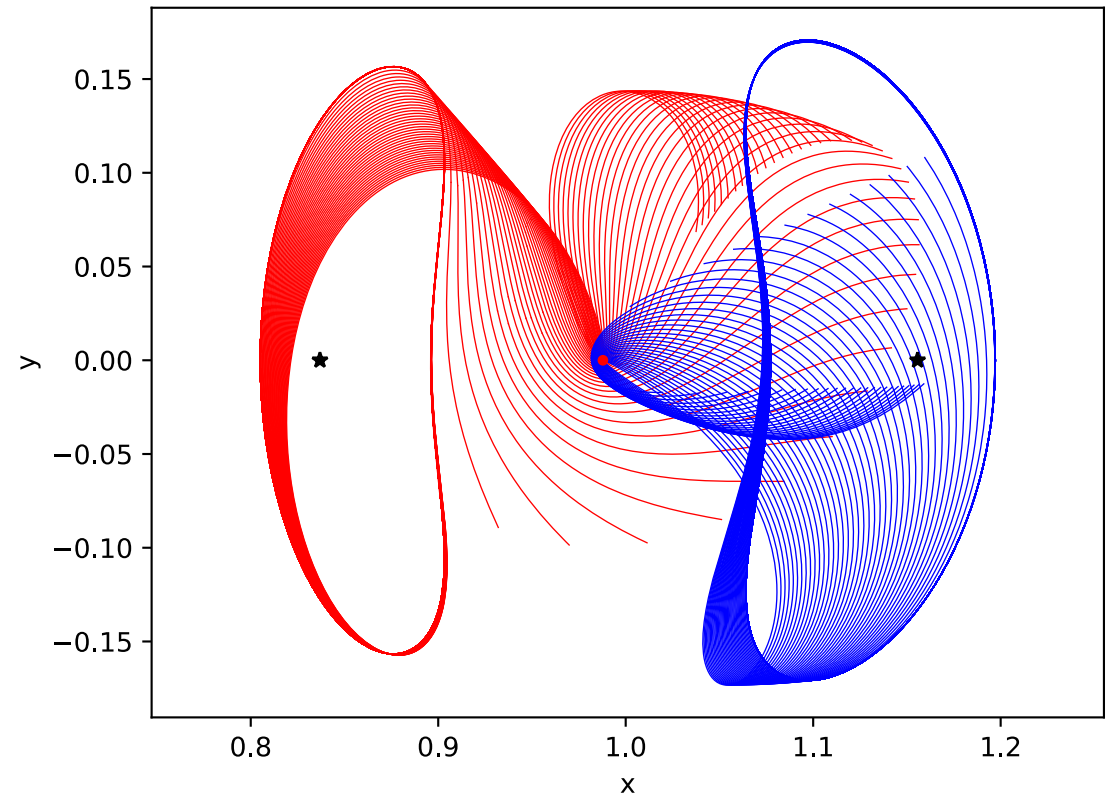
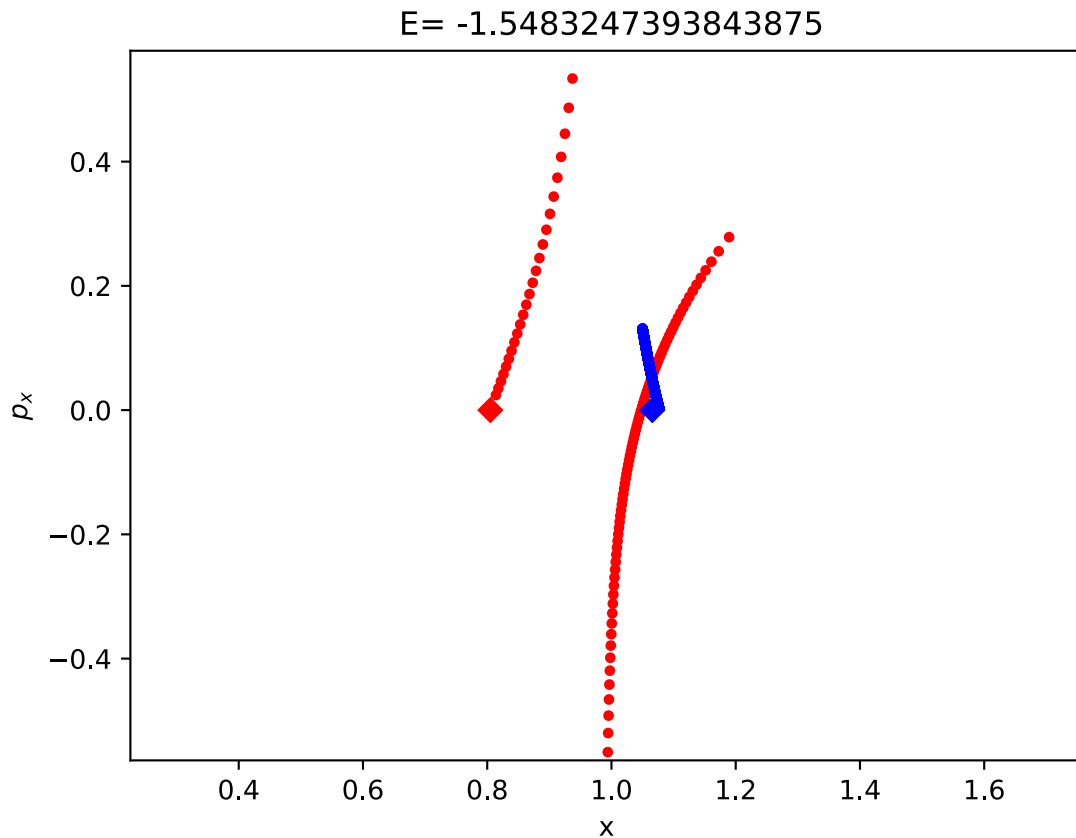
- Pick various points from the manifold in the map
- Integrate backward in time in the flow

- To visualise the unstable tube manifold:

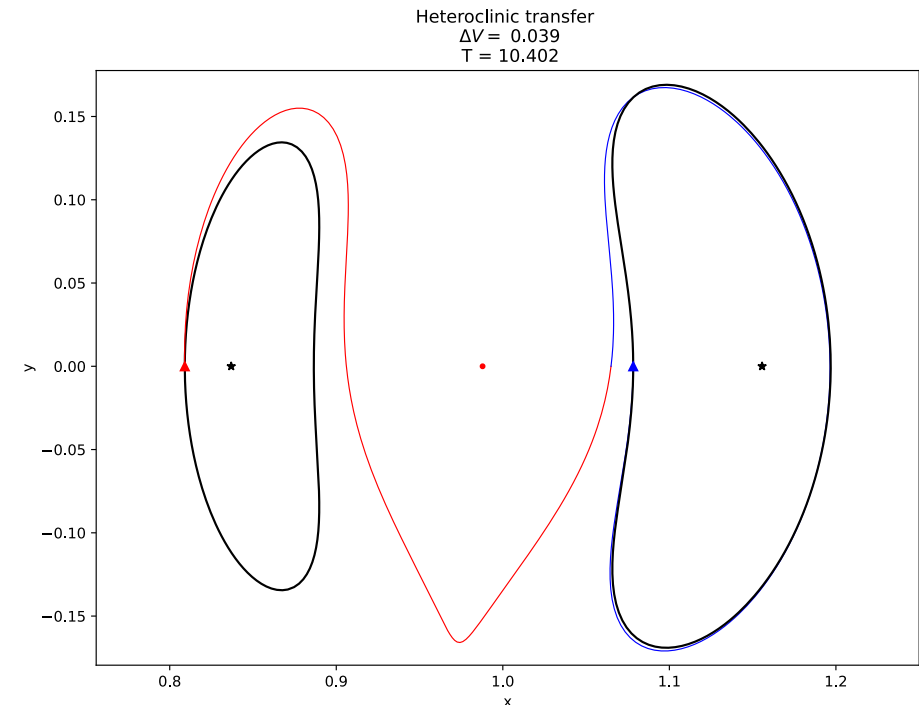
- Pick various points from the manifold in the map
- Integrate forward in time in the flow



- **Heteroclinic intersection:** transverse intersection between the unstable and stable manifolds of two different hyperbolic (or partially hyperbolic) fixed points



- From the intersection integrate backward in time to find the Lyapunov orbit around  $L_1$
- Calculate the  $\Delta V$  for orbit departure
- Transfer
- Arrive in proximity of Lyapunov orbit around  $L_2$
- Possible  $\Delta V$  for orbit insertion
- **Example:**
  - Earth - Moon system,  $\mu = 0.01215$
  - $\Delta V = 0.039$
  - $T = 10.402$



- Lyapunov theorem → Existence of families of closed orbits around collinear L points
- Local manifold theorem → Existence of tube manifolds of horizontal Lyapunov orbit
- Heteroclinic intersections → Low cost transfer trajectories
- Numerical computation:
  - Newton's method applied to the section map  $\Sigma$  in order to compute the periodic orbits
  - Linearization of the Poincaré map  $\Pi$  and use of the local manifold theorem to identify initial conditions on the sets  $W_P^U$   $W_P^S$
  - Numerical propagation of the initial conditions
- Applications:
  - Great distances covered with minimum fuel expenditure
  - Orbits with useful properties (space observatories, communication satellites etc.)
  - Long travel times, trade-off between speed and fuel consumption

- Lecture notes, Dynamical systems; Massimiliano Guzzo; Università degli Studi di Padova
- Dynamical Systems, the Three-Body Problem and Space Mission Design; Koon, Lo, Marsden, Ross; 2011
- The dynamics around the collinear equilibrium points of the RTBP; Gómez, Mondelo; 2001
- Elementary Symplectic Topology and Mechanics; Franco Cardin; Springer